

A Serendipitous XMM Survey of the SDSS: the evolution of the colour-magnitude diagram of X-ray AGN from $z = 0.8$ to $z = 0.1$

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ABSTRACT

A new serendipitous XMM survey in the area of the Sloan Digital Sky Survey is described (XMM/SDSS), which includes features such as the merging of overlapping fields to increase the sensitivity to faint sources, the use of a new parametrisation of the XMM point spread function for the source detection and photometry, the accurate estimation of the survey sensitivity. About 40 000 X-ray point sources are detected over a total area of 122 deg². A subsample of 209 sources detected in the 2–8 keV spectral band with SDSS spectroscopic redshifts in the range $0.03 < z < 0.2$, optical magnitudes $r < 17.77$ mag and $\log L_X(2–10 \text{ keV}) > 41.5$ (erg/s) are selected to explore their distribution on the colour magnitude diagram. This is compared with the colour-magnitude diagram of X-ray AGN in the AEGIS field at $z \approx 0.8$. We find no evidence for evolution of the rest-frame colours of X-ray AGN hosts from $z = 0.8$ to $z = 0.1$. This suggests that the dominant accretion mode of the AGN population, which is expected to imprint on the properties of their host galaxies, does not change since $z = 0.8$. This argues against scenarios which attribute the rapid decline of the accretion power of the Universe with time (1 dex since $z = 0.8$) to changes in the AGN fueling/triggering mode.

Key words: astronomical data bases: surveys – methods: data analysis – galaxies: active – galaxies: Seyferts – X-rays: diffuse background

1 INTRODUCTION

Understanding the evolution of Active Galactic Nuclei (AGN), which signpost accretion events onto Supermassive Black Holes (SMBHs), remains a challenge for modern astrophysics. Although observations have demonstrated beyond any doubt that the luminosity density of these systems has dropped by more than 1 order of magnitude from $z \approx 1$ to the present day (e.g. Ueda et al. 2003; Hasinger et al. 2005; Ebrero et al. 2009; Aird et al. 2010), the physical mechanisms that drive this rapid decline are still not well constrained. The lack of a physical description for the cosmological evolution of SMBHs has implications that go beyond the AGN community. Recent evidence indicates an intimate relation between the building of galaxies and the growth of the SMBH at their centres (e.g. Ferrarese & Merritt 2000; Gebhardt et al. 2000). Therefore without a better understanding of AGN evolution our picture for the buildup of galaxies will also be incomplete.

Galaxy mergers have long been proposed as the mechanism that triggers AGN and drives their cosmological evolution. Numerical SPH (Smoothed Particle Hydrodynamic) simulations demonstrate that these violent events are very efficient in funneling gas to the nuclear galaxy regions (e.g. Hernquist 1989; Barnes & Hernquist 1991, 1996), where it can be consumed by the SMBH (Springel et al. 2005; Di Matteo et al. 2005). Consequently in most

semi-analytic cosmological simulations of galaxy formation (e.g. Cattaneo et al. 2005; Somerville et al. 2008; Wang et al. 2008) the merging of near equal mass gas rich galaxies is the primary mechanism for growing SMBHs. In this family of models the evolution of AGN is intimately related to the decline of the fraction of gaseous major mergers with redshift (e.g. López-Sanjuán et al. 2009). Alternative scenarios propose that the dominant mode of accretion onto SMBHs changes with time (Hasinger 2008; Fanidakis et al. 2010), thereby leading to the observed decline of the AGN space density since $z \approx 1–2$. It is suggested for example, that SMBHs form primarily in violent gaseous major merger events at high redshift ($z \gtrsim 1$; “QSO-mode” Hopkins et al. 2006), while stochastic accretion (e.g. internal instabilities, minor interactions) dominates the growth of SMBHs (“Seyfert-mode” Hopkins & Hernquist 2006) at lower redshift ($z < 0.5$) and produces, on average, lower luminosity systems. Another possibility proposed by Fanidakis et al. (2010) is that the dominant AGN fueling mode shifts from disk instabilities at high redshift to halo gas accretion (also termed “radio” mode, Croton et al. 2006) at low redshift. One of the testable predictions of those models is that the properties of AGN hosts, such as the star-formation history and the morphology, change with redshift, from $z \gtrsim 1$ to $z < 0.5$. In contrast, in the merger only

2 Georgakis & Nandra

driven evolution scenario, AGN should be hosted by galaxies that have similar properties at all redshifts.

The study of AGN hosts as a function of redshift requires a homogeneously selected AGN sample over a wide redshift baseline, $z \approx 0$ to $z \approx 1$ and beyond. Only then can one compare directly the properties of AGN hosts at different epochs. Variations in the selection function with redshift are hard to account for and may lead to erroneous conclusions. At $z \gtrsim 1$ in particular, X-ray observations, especially at energies > 2 keV, are one of the most efficient and least biased methods for locating AGN with a selection function that is easy to quantify. The infrared is a promising wavelength regime for finding active SMBHs, although there are issues related to contamination of infrared selected AGN samples by starbursts (e.g. Georgantopoulos et al. 2008; Donley et al. 2008; Pope et al. 2008; Georgakakis et al. 2010). Optical spectroscopy is also a powerful tool for identifying AGN, but aperture effects, which are particularly severe at high redshifts, raise concerns about dilution of the AGN emission lines by the host galaxy stellar population (e.g. Severgnini et al. 2003). As a result, surveys with XMM and Chandra, both deep/pencil-beam and shallow/wide, have been the workhorse of the astronomy community for compiling AGN samples. Those surveys however, essentially probe active SMBHs close to the peak of the accretion power of the Universe, $z \gtrsim 1$, and lack the area to provide meaningful constraints on the AGN population at $z \lesssim 0.5$. The Sloan Digital Sky Survey (SDSS; Abazajian et al. 2009) has identified the largest sample of low redshift ($z \approx 0.1$) AGN to date, using diagnostic emission line ratios (Kauffmann et al. 2004). The selection function of that sample however, is very different from that of X-ray AGN in deep surveys, rendering the comparison difficult. For example, the SDSS AGN include a large number of LINERs (Low Ionisation Nuclear Emission line Region), which are controversial objects and may not be powered by accretion onto a SMBH (e.g. Sarzi et al. 2010). More relevant to high redshift X-ray AGN surveys are the serendipitous near all-sky AGN samples compiled in the nearby Universe by the high energy missions INTEGRAL (20–100 keV, total of 144 Seyfert AGN, Beckmann et al. 2009) and SWIFT (15–195 keV, total of 266 Seyferts, Tueller et al. 2010). Although those AGN samples are selected at much higher rest-frame energies compared to X-ray sources detected by Chandra or XMM at $z \approx 1$ (< 15 keV), they are more appropriate than SDSS AGN as a low redshift comparison sample. Ideally however, one would like to select nearby AGN at rest-frame X-ray energies comparable to those of Chandra/XMM surveys at $z \approx 1$.

In this paper we describe the compilation of a large sample of X-ray selected AGN at $z \approx 0.1$ using a new serendipitous X-ray survey (hereafter referred to as XMM/SDSS) in the area of the SDSS based on archival XMM observations. The advantage of the low redshift X-ray AGN subset of the XMM/SDSS survey is that the selection function is almost identical to deep pencil-beam samples, thereby minimising differential selection biases. As a result the comparison between the AGN host galaxy properties between $z \approx 0.1$ and $z \approx 1$ is greatly facilitated. The XMM archive includes over 10 years worth of observations, which allow serendipitous surveys over many tens of square degrees on the sky. Therefore the volume probed by the XMM/SDSS sample at $z \approx 0.1$ is orders of magnitude larger than any wide-area survey carried out by either XMM or Chandra, thereby allowing detailed studies of the statistical properties of low redshift X-ray AGN. The SDSS is the field of choice for our serendipitous X-ray survey as it provides the essential follow-up observations for the identification of low redshift X-ray AGN. The SDSS optical data include 5-band photometry (*ugriz*) and extensive optical spectroscopy for galaxies at $z \approx 0.1$.

Wide area surveys at various wavelengths in the SDSS are either underway or completed, thereby allowing panchromatic studies of the XMM/SDSS X-ray AGN over a wide redshift baseline. The ancillary data in the SDSS include (i) the UKIRT Infrared Deep Sky Survey (UKIDSS, Lawrence et al. 2007), which will cover 4000 deg^2 of the SDSS in 4 near-infrared bands to $K = 18.5$ mag, (ii) the AKARI all-sky survey in 6 mid- and far-infrared bands from $9 - 200 \mu\text{m}$ (Ishihara et al. 2010), (iii) the Herschel ATLAS far-infrared survey (Eales et al. 2010) and (iv) the FIRST survey, which provides deep (1 mJy) radio continuum (1.4 GHz) images of the entire SDSS (Becker et al. 1995).

The structure of the paper is as follows. In sections 2 to 9 we describe the automated pipeline developed to reduce the archival XMM observations that overlap with the SDSS. The data reduction steps include the construction of event files and images in different energy bands, the detection of sources, the calculation and application of astrometric corrections, the estimation of fluxes and the identification of X-ray sources with optical counterparts. The pipeline uses a new parametrisation of the XMM Point Spread Function (Appendix A), which is employed for source detection and photometry, and an accurate method for estimating the sensitivity of the X-ray survey to point sources. Section 10 demonstrates the utility of our serendipitous X-ray AGN sample by comparing, for the first time, the colour magnitude diagrams of X-ray selected AGN at $z \approx 0.1$ and $z \approx 0.8$. Section 11 discusses the results and conclusions. Throughout this paper we adopt $H_0 = 100 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.3$ and $\Omega_\Lambda = 0.7$. Rest frame quantities (e.g. absolute magnitudes, luminosities) are parametrised by $h = H_0/100$.

2 ORGANISATION OF THE XMM DATA

For the construction of the serendipitous X-ray source catalogue, XMM observations that were taken in full frame and extended full frame modes and were public by July 2009 (total of 3818 observations) are selected. We then chose only those that overlap with the SDSS DR7 (Data Release 7; Abazajian et al. 2009). We use the spherical polygons (Hamilton & Tegmark 2004) defined in the New York University Value Added Galaxy Catalogue (NYU-VAGC; Blanton et al. 2005) to describe the geometry of the SDSS survey. This geometry is the area covered by the SDSS-DR7 imaging survey after masking out Tycho stars as described by Blanton et al. (2005). If the aimpoint of an XMM observation falls within one of the masked areas it is excluded from further analysis. This requirement reduced the number of analysed XMM observations to 928.

The basic unit of the survey data structure is the individual XMM observation, which will also be referred to as “obsid”. XMM observations are assigned to “projects” and “pointings”, which are defined below. Overlapping obsids are merged by the pipeline to maximise the sensitivity of the survey to faint sources. For that purpose the data are organised into projects which include at least one XMM observation. If an XMM field lies within 28 arcmin of the aimpoint of at least one obsid of a project, it is assigned to that project. Otherwise a new project is defined. The radial distance cut is chosen to ensure overlap between adjacent XMM fields, which have usable sizes of approximately 14 arcmin radius. Observations that belong to separate projects are therefore non overlapping. A project is assigned a unique name using up to the first 9 characters of the target name keyword of the first XMM observation assigned to that project. Projects with more than one XMM observations are

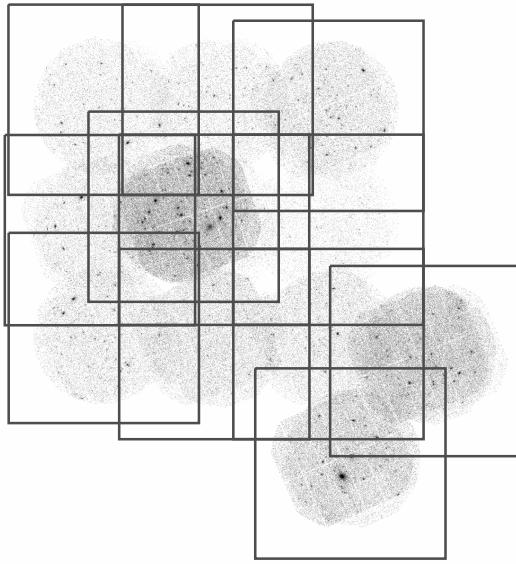


Figure 1. Demonstration of how the XMM observations are organised into projects and pointings. The individual XMM observations shown in the figure belong to the same project as they are contiguous. The identification name of the project is SDSS1341 after the target name of the first obsid that was assigned to the project. The squares are 40 arcmin on the side and represent pointings. All XMM observations that overlap with a 40×40 arcmin 2 box are merged to increase the sensitivity of the serendipitous survey to faint sources. The positions of the pointings are defined in a way which guarantees that the entire contiguous area of the project is covered by at least one pointing independent of the layout of the individual obsids

organised in pointings, which are defined as 40×40 arcmin 2 boxes and include all the obsids of the project that overlap with that box. The first observation of a project defines the centre of the first pointing. Project observations that lie more than 6 arcmin away from an existing pointing define the centres of new pointings. This choice of radial distance leads to substantial overlap between adjacent pointings but ensures that the contiguous area of a project is covered by at least one pointing independent of the layout of individual obsids. This is demonstrated in Figure 1 which plots the positions of the pointings of a particular project in the XMM serendipitous survey. Once the pointings of a project are defined, they are assigned all the project observations that lie within 28 arcmin off the centre of each pointing. Therefore all the XMM observations that have at least a small fraction of their field of view overlapping with the 40×40 arcmin 2 pointing box belong to the same group and are eventually merged by the pipeline. A particular obsid of a project may belong to more than one pointings.

3 DATA PROCESSING

The XMM observations are reduced using the Science Analysis System (SAS) version 9.0. The first step is to produce event files from the Observation Data Files (ODF) using the EPCHAIN and EMCHAIN SAS tasks for the EPIC PN and MOS detectors respectively. This includes (i) the creation of raw event lists and Good Time Interval (GTI) data for each CCD of the EPIC detectors, (ii) the identification of bad pixels, (iii) the calibration of the raw event lists by flagging trailing events, performing pattern recognition, gain and CTI (Charge Transfer Inefficiency) corrections, (iv) the

estimation of linearised detector and sky coordinates and (v) the merging of the calibrated event lists of different CCDs into event files for each EPIC detector.

Pixels along the edges of the CCDs of the PN and MOS detectors are removed, as in our experience their inclusion often results to spurious detections. The excluded pixels for the EPIC-PN correspond to rows with RAWX = 1, 64 and columns with RAWY = 199 – 200 and RAWY < 24. The latter set of pixels are flagged out to reduce bright low energy edges (Watson et al. 2009). For the EPIC MOS1, MOS2 we remove pixels with coordinates (RAWX, RAWY) = (1–4, 1–600), (RAWX, RAWY) = (597 – 600, 1 – 600), (RAWX, RAWY) = (1 – 600, 1 – 4), (RAWX, RAWY) = (1 – 600, 597 – 600). Events with patterns > 4 for the PN and > 12 for the MOS detectors are removed and the filter formulae (FLAG & 0x766b0808)==0 for the MOS and (FLAG & 0x2fb0808)==0 for the PN are applied. Those flags exclude events from offset columns and spoiled frames and remove electronic noise (Marty et al. 2003). Further filtering is applied to observations targeting bright X-ray sources, such as QSOs or Galactic stars (total of 47). Those observations are identified by visual inspection and are found to suffer from diffraction spikes as well as out-of-time events stripes. The targets are placed at the aimpoint of the XMM EPIC cameras and therefore the affected regions are fixed in detector coordinates. For the PN we remove the out-of-time events stripe in rows RAWX = 29 – 44 of CCD4 and the upper part of CCDs 4 (RAWY > 155) and 7 (RAWY > 170), which are affected by the diffraction spikes of the target. In the case of MOS1 the removed pixels lie on CCD 1 (RAWX = 292 – 328 and RAWY = 1 – 600, RAWX = 140 – 480 and RAWY = 164 – 454). For MOS2 we remove pixels from CCD 1 with detector coordinates RAWX = 277 – 320 and RAWY = 1 – 600, RAWX = 163 – 447 and RAWY = 144 – 444.

Flaring background periods are identified using a methodology similar to that described in Nandra et al. (2007). First, sources are removed from the event file to avoid AGN variability biasing the results. The EWAVELET task of SAS with a threshold of 5 times the local background rms is used to detect sources in the 0.5–8 keV spectral band of the combined image of all EPIC detectors available to a particular XMM observation. These sources are then masked out from the event file before generating the background light curve by binning events in the energy range 0.2–12 keV in 20 s intervals. The approximate quiescent background level is determined by calculating the count rate limit at which the excess variance of the background light curve is minimum (c.f. Nandra et al. 2007, who estimates the limit where the excess variance is zero). Periods with background count rate 2 times higher than the level where the excess variance is minimum are excluded from the filtered event file (see Figure 2). It is verified that this methodology works well for XMM. It fails to identify however, cases where the entire observation is affected by flares and hence, has an overall elevated background. Therefore PN and MOS observations with quiescent background levels higher than 50 and 20 ct s $^{-1}$ respectively, are not analysed. It is empirically found that those values exclude problematic observations from further analysis. Also, if the Good Time Interval for a particular detector, after the removal of flaring periods, is less than 1 ks, it is excluded from further analysis. A total of 95 obsids have been removed entirely either because of high flares or short clean exposure times.

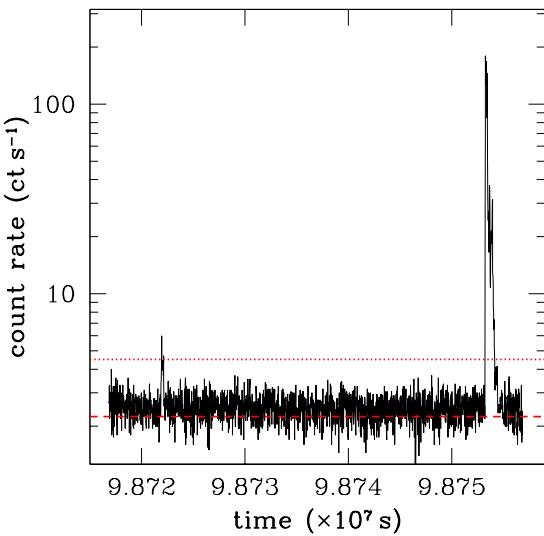


Figure 2. Light curve of XMM observation 0041170101 which shows a flare at the end of the observation. The dashed horizontal line marks the approximate quiescent background level determined by the methodology described in the text. The dotted line corresponds to the count rate limit above which events are excluded.

4 IMAGE CONSTRUCTION

Images in celestial coordinates with pixel size of 4.35 arcsec are constructed in 5 energy bands, 0.5-8 keV (full), 0.5-2 keV (soft), 2-8 keV (hard), 5-8 keV (very hard; vhrd) and 7.5-12 keV (ultra hard; uhrd). The upper energy limit of 8 keV for the full, hard and vhrd band images is chosen because of the reduced effective area of XMM at higher energies. Also, below 0.5 keV the XMM background is elevated due to the Galactic X-ray emission component (Lumb et al. 2002). Therefore, photons with energies < 0.5 and > 8 keV primarily increase the background and do not improve the signal-to-noise ratio of the final image, thereby limiting the sensitivity of the survey. Exposure maps, which account for vignetting, CCD gaps and bad pixels, are constructed at monochromatic energies of 2 keV (full band), 1 keV (soft), 3.5 keV (hard), 6 keV (vhrd) and 8.5 keV (uhrd). Those values correspond to the mean photon energy in each spectral band of a power-law X-ray spectrum with $\Gamma = 1.4$ (i.e. similar to the diffuse X-ray background; e.g. Hickox & Markevitch 2006) folded through the sensitivity of the XMM EPIC cameras (mirror effective area, energy redistribution matrix).

For projects with at least one pointing, the EPIC images and exposure maps of all the observations of a pointing are merged prior to source detection to maximise the sensitivity to faint sources. In the case of projects that consist of a single XMM observation the EPIC images and exposure maps are coadded.

5 POINT SOURCE DETECTION

X-ray point sources are searched for in the merged EPIC images of either a pointing or a project with a single XMM observation. The detection of sources is performed independently in each of the 5 spectral bands defined in section 4, soft, full, hard, vhrd and uhrd. A methodology similar to that described in Nandra et al. (2005) and Laird et al. (2009) is adopted. Source candidates are identified using the wavelet based Ewavelet source detection task of SAS at

a low threshold of 4σ above the background, where σ is the rms of the background counts.

For each candidate source the Poisson probability of a random background fluctuation is estimated. This step involved the extraction of the total counts at the position of the source and the determination of the local background value. For the extraction of counts a new set of Point Spread Function (PSF) estimates is employed, which accounts for both the varying elliptical shape and size of the XMM PSF as a function of position on the EPIC detectors. Details of the calculation of the PSF are presented in the appendix A. At each source position the counts are extracted within an elliptical aperture that corresponds to 70 per cent of the PSF Encircled Energy Fraction (EEF). The total counts at a candidate source position, T , is the sum of the extracted counts from individual EPIC cameras. For each source the local background is estimated by first masking out all detections within 4 arcmin of the source position using an elliptical aperture that corresponds to the 80 per cent EEF ellipse. The counts from individual EPIC cameras are then extracted using elliptical annuli centred on the source with inner and outer semi-major axes of 5 and 15 pixels (0.36 and 1.09 arcmin) respectively. The position angle and ellipticity of the ellipses are set equal to those of the 70 per cent EEF. The mean local background, B , is then estimated by summing up the background counts from individual EPIC cameras after scaling them down to the area of the source count extraction region. The Poisson probability $P(T, B)$ that the extracted counts at the source position, T , are a random fluctuation of the background is calculated. We consider as sources the detections with $P(T, B) < 4 \times 10^{-6}$. At this cutoff less than about 0.01 false sources per XMM pointing are expected given the typical size of the 70 per cent EEF aperture (15-20 arcsec radius depending on energy and off-axis angle) and the field of view of the XMM (about 14 arcmin radius). The candidate source list generated by Ewavelet task includes a large number of background fluctuations which are masked out during the estimation of the local background for individual sources. This may lead to an underestimation of the local background, as positive background fluctuations may be excluded as source candidates. We account for this potential source of bias by repeating the process above using as input candidate source list the one with Poisson probability $P(T, B) < 4 \times 10^{-6}$.

Our methodology is optimised for the detection of point sources. The final catalogue however, includes extended X-ray sources associated with hot gas from galaxy clusters or groups. Also, the extended X-ray emission regions of bright clusters are often split into multiple spurious detections by our source detection pipeline. The SAS task EMLDETECT, although also optimised for the detection of point sources, can also identify extended X-ray emission (see Watson et al. 2009 for a discussion of limitations) as well as spurious detections. This task is used to perform simultaneous PSF fits to the source count distribution of all the EPIC images of an XMM observation to estimate the detection likelihood, the source extent and the extent likelihood. EMLDETECT is applied separately to each of the five energy bands, soft, full, hard, vhrd and uhrd. The free parameters are the source count rate and source extent. The source positions are kept fixed during the minimisation process. The background maps described in section 7 are used as input to the EMLDETECT task. The energy and position dependent PSF provided in the calibration database is used in the fit. We caution that the PSFs employed by EMLDETECT are different from those described in Appendix A. The task is allowed to fit up to 3 sources simultaneously if they fall within the 90 per cent EEF radius of the calibration database PSF. For the determination of the

extent of a source the PSF is convolved with a β -model profile. If the likelihood of the extended model is below 3 or the extent of the source is less than 1.5 pixels then point source parameters are determined by the task. In the case of pointings with multiple observations the EMLDETECT task is ran on each observation separately. The individual likelihoods for a given source are then summed up and transformed to equivalent likelihoods L_2 corresponding to the case of two free parameters, as described in the EMLDETECT documentation. Based on visual inspection of the XMM images we consider as extended those sources with extension likelihood > 3 in either the full or the soft spectral bands. For point sources we wish to keep the sample selection as clean and transparent as possible. One of the advantages of our point-source detection methodology is that the selection function, which is based on the Poisson spurious probability threshold, can be quantified to a high degree of accuracy (see sensitivity map section below). Introducing additional selection criteria, e.g. keeping sources above a given EMLDETECT detection likelihood threshold, would modify the selection function of the sample in a way that is hard to quantify. We therefore choose to keep those biases to the minimum by excluding only those point sources for which EMLDETECT failed to determine a reliable fit and did not estimate a detection likelihood. In the final XMM/SDSS catalogue those sources are listed with an EMLDETECT detection likelihood < 0 and are considered spurious. Such sources are typically associated with image artifacts, such as diffraction spikes that have not been properly removed, extended hot gas emission split by the detection algorithm into multiple sources etc.

The source catalogues in different energy bands are merged to produce a unique list of sources for either a pointing or a project consisting of a single XMM observation. A source detected in a particular energy band is cross-matched with sources in other energy bands by finding the closest counterpart. The search radius is set to the 70 per cent EEF circular radius, i.e. $SMA \times \sqrt{1 - \epsilon}$, where SMA is the semi-major axis and ϵ is the ellipticity of the ellipse. If a source lies on multiple obsids or detectors the search radius is estimated as the exposure time weighted average of individual 70 per cent EEF circular radii. It is verified that this choice of search radius is the optimal for identifying the same source in different energy bands. In practice we start by band-merging the full band sources and then proceed with the soft, hard, vhrd and uhrd catalogues.

As there is overlap between adjacent pointings of a project, it is essential to identify and remove duplicates to produce unique source lists for projects with multiple pointings. For each source of a pointing we search for the closest counterpart in adjacent pointings within a fixed search radius of 12 arcsec. From the list of duplicates we keep in the final list the one with the smaller angular distance from the aimpoint of the pointing on which it is detected.

In the final band-merged source catalogue of unique sources there are 1258 extended and 827 spurious sources and a total of 39830 point sources. Table 1 presents the total number of point sources detected in each spectral band. It should be stressed that our detection method is geared toward point sources and therefore the extended source catalogue should be treated with caution. Providing a well defined sample of X-ray selected clusters is *not* among our goals. The final point source catalogue also includes targets.

6 ASTROMETRY

The pointing errors of the XMM observations are of the order of few arcsec. The positions of individual X-ray sources how-

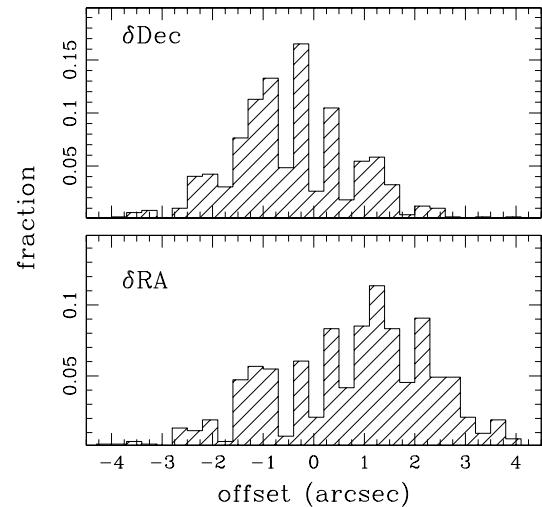


Figure 3. Distribution of the RA (bottom panel) and DEC (top panel) corrections applied to the celestial coordinates of the X-ray sources to account for systematic errors in the determination of the absolute XMM pointing. The δRA and δDEC offsets are estimated by matching the positions of X-ray sources with optical counterparts in the SDSS.

ever, can be determined to an accuracy of 1-2 arcsec (Waston et al. 2009). The astrometry of the X-ray catalogue can therefore be improved by correcting the RA and Dec of X-ray sources for systematic errors. For that purpose the positions of optically selected sources in the SDSS, which have sub-arcsec astrometric uncertainties ($\lesssim 0.1$ arcsec, Pier et al. 2003), are cross-correlated with the X-ray source list. The EPOSCORR task of SAS is used to determine the offsets in RA, Dec and roll angle of an XMM pointing that maximises the number of optical identifications. Only X-ray point sources are used in this exercise, i.e. detections flagged as extended or spurious by the EMLDETECT task of SAS are excluded from the cross-matching. Also, only SDSS sources brighter than $r = 22$ mag are matched to the X-ray source positions. In the case of observations with less than 8 optical identifications, no corrections are applied to the X-ray source positions. Figure 3 plots the distribution of the RA, Dec offsets applied to the celestial coordinates of the X-ray sources and shows that typical corrections are about 1-2 arcsec. In the case of projects with multiple pointings the astrometry of individual obsids should be corrected for systematics prior to merging. We have verified however, that this approach produces results that are similar to the simpler and less time-consuming method of estimating astrometric corrections for the final source list of individual pointings of a project, i.e. after merging.

7 FLUX ESTIMATION AND HARDNESS RATIOS

For sources with Poisson detection probability $< 4 \times 10^{-6}$ in at least one of the five spectral bands fluxes are estimated in all bands by summing up the counts within the 80 per cent EEF of the elliptical PSFs described in Appendix A. Compared to source detection, a larger aperture is adopted for the photometry. The choice of radius is to include as large a fraction of source photons as possible, while at the same time keeping the background noise low. Also, for larger EEFs, the radius becomes large enough that the contribution of pho-

6 Georgakis & Nandra

tons from nearby sources becomes an issue. The local background at the position of a source is estimated as described in the previous section. The counts from all EPIC cameras of all the XMM observations that a source lies on are summed up for the estimation of the flux in a particular energy band. The different response of the EPIC cameras is accounted for by adopting the appropriate count to flux conversion factors, ECF. Suppose a source on the detector i of an XMM observation. The total number of counts, C_i , at the source position within the 80 per cent EEF are

$$C_i = f_X \times ECF_i \times EEF \times t_i + B_i, \quad (1)$$

where f_X is the source flux, ECF_i is the count to flux conversion factor for the camera i , $EEF = 0.8$ is the encircled energy fraction, t_i is the exposure time at the source position and B_i is the local background value. Summing up over i , i.e. over all observations and all EPIC cameras on which the source lies on gives

$$\sum_i C_i = f_X \times EEF \times \sum_i (ECF_i \times t_i) + \sum_i B_i. \quad (2)$$

The simplest approach for estimating the source flux is to solve the equation above for f_X

$$f_X = \frac{1}{EEF} \sum_i (C_i - B_i) / \sum_i (ECF_i \times t_i). \quad (3)$$

The equation above however, ignores the Eddington bias and the fact that there is a continuous probability distribution of fluxes which can produce the observed counts. We account for these effects by adopting the Bayesian methodology described in Laird et al. (2009) to calculate the source fluxes and 1σ confidence limits from the detected counts. In brief, the probability of the source having flux f_X given the observed number of total counts C (source and background) follows Poisson statistics

$$P(f_X, C) = \frac{T^C e^{-T}}{C!} \pi(f_X) \quad (4)$$

where T is the mean expected total counts in the detection cell. This quantity is determined as a function of flux from equation 1. The last term in equation 4, $\pi(f_X)$, reflects our prior knowledge of the distribution of source fluxes, i.e. the differential X-ray source number counts, which are described by a double power-law (e.g. Georgakis et al. 2008). This term accounts for the Eddington bias, in which statistical variations of the observed source counts combined with the steep $\log N - \log S$ of the X-ray population result in brighter measured fluxes for the detected sources compared to their intrinsic ones. The differential X-ray source counts are approximated by a double power-law with faint and bright end slopes of -1.5 and -2.5 respectively and a break flux of 10^{-14} erg s $^{-1}$ cm $^{-2}$. The adopted values are representative of the double power-law parameters estimated for X-ray samples selected in different energy bands (e.g. Georgakis et al. 2008). Equation 4 is used to estimate the mode of the X-ray flux distribution, which can be calculated analytically. The 68 per cent confidence limits on the flux are estimated by integrating equation 4 as described in Laird et al. (2009).

The energy to flux conversion factor is estimated for each EPIC camera separately assuming a power-law X-ray spectrum with $\Gamma = 1.4$, i.e. similar to the XRB, absorbed by the appropriate Galactic hydrogen column density. The latter is derived from the HI map of Kalberla et al. (2005) using the right ascension and declination of the aimpoint of each XMM observation and the NH task of FTOOLS. The energy to flux conversion factors are such that the counts from the soft, full, hard, vhrd and uhrd bands are trans-

formed to fluxes in the 0.5-2, 0.5-10, 2-10, 5-10 and 7.5-12 keV bands respectively.

Hardness ratios (HRs) are estimated between the soft (0.5-2 keV) and the hard (2-8 keV) band count rates, S and H respectively, following the standard definition

$$HR = \frac{H - S}{H + S}. \quad (5)$$

The BEHR (Bayesian Estimation of Hardness Ratios, Park et al. 2006) code is used, which is designed to determine HRs in the Poisson regime of low counts and to compute the proper uncertainty regardless of whether the source is detected in both passbands or not. Hardness ratios were estimated independently for the PN, MOS1 and MOS2 detectors. In the case of multiple overlapping observations the summed source and background counts as well as the total exposure time from each detector are used as input to the BEHR code. The estimated HRs are scaled to the on-axis exposure time by setting the BEHR parameters SOFTEFF and HARDEFF to the ratios of the on-axis exposure time over the exposure time at the source position in the soft and hard bands respectively. In the case of overlapping observations these parameters are defined by the relation

$$eff = \frac{\sum_i C_i \frac{t_{on,i}}{t_i}}{\sum_i C_i}, \quad (6)$$

where eff represents either the SOFTEFF or the HARDEFF BEHR parameter, the summation is over all XMM observations, i , C_i represents the counts (soft or hard band) at the source position (source and background), $t_{on,i}$ is the on-axis exposure time and t_i is the exposure time at the source position.

8 SENSITIVITY MAPS

The construction of the sensitivity maps follows the methodology of Georgakis et al. (2008), which accurately estimates the probability of detecting a source with a given X-ray flux accounting for observational effects, such as vignetting, flux estimation biases (e.g. Eddington bias) and the fraction of spurious sources expected in any source catalogue. The important parameters in this exercise are the size and shape of the detection cell, which are well defined in our source detection method, and the Poisson probability threshold, P_{thresh} , below which an excess of counts is considered a source. By fixing P_{thresh} one sets the minimum number of photons in a cell, L , for a formal detection.

The first step is the estimation of the source-free background across the EPIC detectors. The background map is estimated using custom routines to first remove the counts in the vicinity of detected sources using the 80 per cent EEF elliptical apertures. Pixel values in the source regions are replaced by sampling from the distribution of pixel values in the local background regions. These are defined by elliptical annuli centred on each source with inner semi-major axis 5 pixels (≈ 21.8 arcsec) larger than the 80 per cent EEF size and widths of 15 pixels (≈ 62.3 arcsec). The resulting maps are smoothed by replacing each pixel value with the median within a sliding box of 20 pixels (87 arcsec) in size. The smoothed background maps are then used to estimate the mean expected background counts, within the detection cell, B (i.e. the 70 per cent EEF elliptical apertures). In the case of merged images from different detectors and/or XMM observations, B is the sum of the mean expected background counts within the detection cell of individual images. The cumulative probability that the observed counts in a particular detection cell will exceed L is

$$P_B(\geq L) = \gamma(L, B), \quad (7)$$

where the function $\gamma(a, x)$ is the incomplete gamma function defined as

$$\gamma(a, x) = \frac{1}{\Gamma(a)} \int_0^x e^{-t} t^{a-1} dt. \quad (8)$$

Adopting a detection threshold P_{thresh} one can invert equation 7 numerically to estimate the (integer) detection limit L for a cell with mean expected background B . Repeating this exercise for different cells across the image one can determine L as a function of position (x, y) on the detector. This 2-D image of L values is the sensitivity map. Note that the sensitivity map is independent of the spectral shape of the source.

A useful 1-D representation of this image, with a wide range of applications, is the total detector area in which a source with flux f_X can be detected. The cumulative distribution of the area plotted as function of f_X is often referred to as sensitivity curve. This was constructed adopting a Bayesian approach as follows. For a source with flux f_X and a given spectral shape ($\Gamma = 1.4$ in this paper) we can determine the probability of detection in a cell with detection limit L and mean background B . The total observed counts in the cell are $C = S + B$, where S is the mean expected source contribution. C is given by equation 2 in the generic case of merged images from different EPIC detectors and/or XMM observations.

Both B and S fluctuate and therefore using equation 7 the probability their sum exceeds the detection threshold is $P_{C, f_X}(\geq L) = \gamma(L, C)$. The sensitivity curve is the sum of the $P_{C, f_X}(\geq L)$ distributions of individual detection cells. The sensitivity curve of the XMM/SDSS survey in different bands is plotted in Figure 4. That figure also shows that the total XMM/SDSS surveyed area is 122 deg^2 . In addition to the Bayesian approach we also estimate sensitivity curves adopting the standard methodology of assigning a single limiting flux to a detection cell, i.e. assuming that the minimum net counts for a source to be detected in the cell are $S = L - B$. Figure 5 compares the Bayesian and standard sensitivity curves in different spectral bands and shows that they differ substantially at the faint flux end.

9 OPTICAL IDENTIFICATIONS

The Likelihood Ratio method (LR; Sutherland & Saunders 1992; Laird et al. 2009) is adopted to identify X-ray sources with optical counterparts in the SDSS. In this exercise the quantities of interest are (i) the probability that an optical source, at a given distance from the X-ray position and with a given optical magnitude, is the true counterpart and (ii) the probability that the same source is a spurious alignment. The likelihood ratio is defined as the ratio between the two probabilities

$$\text{LR} = \frac{q(m) f(r)}{n(m)}, \quad (9)$$

where $q(m)$ is the expected magnitude distribution of the true optical counterparts, $f(r)$ is the probability distribution function of the positional uncertainties in both the X-ray and the optical source catalogues and $n(m)$ is the background density of optical galaxies of magnitude m in the SDSS r -band.

For the positional accuracy of the X-ray sources we adopt a Gaussian distribution with standard deviation of 1.5 arcsec (see below). The a priori probability $q(m)$ that an X-ray source has a counterpart with magnitude m , is determined as follows. Firstly, the

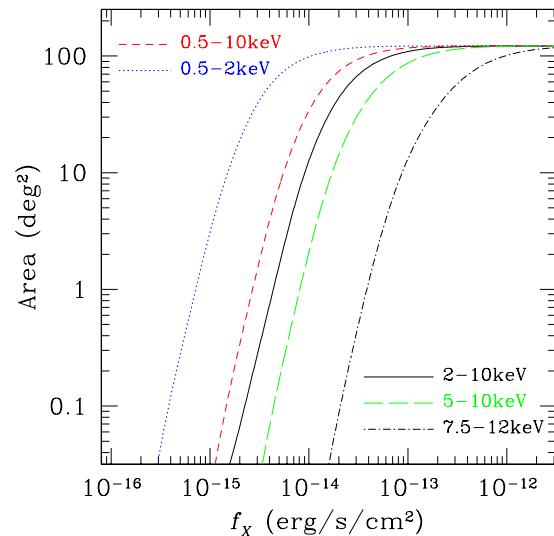


Figure 4. Sensitivity curves of the XMM/SDSS survey in different spectral bands. The curves are red dashed: full-band; blue dotted: soft-band; black continuous: hard-band; green long-dashed: very-hard band; black dashed-dotted: ultra-hard band. The total surveyed area is 122 deg^2 .

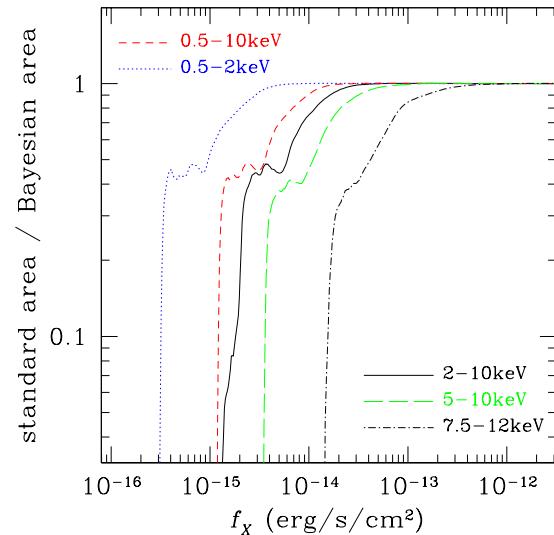


Figure 5. Ratio between the standard and Bayesian sensitivity curve estimates as a function of flux in different spectral bands. The curves are red dashed: full-band; blue dotted: soft-band; black continuous: hard-band; green long-dashed: very-hard band; black dashed-dotted: ultra-hard band. At faint fluxes the standard and Bayesian approaches for determining the sensitivity curve differ substantially.

X-ray and optical source catalogues are matched by simply finding the closest counterpart within a fixed search radius of 4 arcsec. The magnitude distribution of the spurious matches is estimated by scaling $n(m)$ to the area of 4 arcsec radius within which we search for counterparts. This is then subtracted from the magnitude distribution of the counterparts in the real catalogs to determine the magnitude distribution of the true associations, $q(m)$.

Having estimated $f(r)$ and $q(m)$ we identify all possible counterparts to X-ray sources within a 4 arcsec radius and the LR of each one is determined using equation 9. We consider as counterparts those galaxies with LR above a certain limit. The choice of the cutoff in LR is a trade off between maximum number of counterparts and minimum spurious identification rate. In order to assess how secure an optical counterpart is we use the reliability parameter defined by Sutherland & Saunders (1992)

$$\text{Rel}_i = \frac{\text{LR}_i}{\sum_j \text{LR}_j + (1 - Q)}, \quad (10)$$

where Rel_i is the reliability of the i optical counterpart of an X-ray source, the index j of the summation runs over all the possible counterparts within the search radius and Q is the fraction of X-ray sources with identifications to the magnitude limit of the optical survey. It can be shown that the sum of the reliabilities of individual counterparts equals the expected number of true associations (Sutherland & Saunders 1992). Comparison of $\sum_i \text{Rel}_i$ with the total number of counterparts with $\text{LR} > \text{LR}_{\text{limit}}$ provides an estimate of the spurious identification rate. By varying LR_{limit} one can minimise the number of false associations. We chose to use $\text{LR}_{\text{limit}} = 1.5$. At this cutoff 49 per cent (19431/39830) of the X-ray point sources have counterparts with an estimated spurious identification rate of about 7 per cent. Table 1 presents details on the total number of optical identifications for different spectral bands as well as the number of sources in subsamples defined by applying various selection criteria, e.g. magnitude cuts, spectroscopic redshift availability etc.

The optical identification fraction for the full band is plotted as function of 0.5–10 keV flux limit in Figure 6. The optical identification fraction increases to about 90 per cent at bright fluxes, $f_X \approx 10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2}$. The magnitude distribution of the optical counterparts of the XMM/SDSS sources is shown in Figure 7. The distribution of X-ray sources in the X-ray to optical flux ratio (f_X/f_{opt}) plane is shown in Figure 8. A total of 2067 X-ray point sources have spectroscopic observations from the SDSS. The redshift histogram is shown in Figure 9. For X-ray sources associated with galaxies (i.e. optically resolved) brighter than $r < 17.77$ mag (the magnitude cut of the SDSS main galaxy spectroscopic sample) the redshift sampling rate is about 90 per cent, i.e. similar to that of the overall SDSS galaxy population with $r < 17.77$ mag. The majority of spectroscopically unidentified X-ray sources that are brighter than this limit are associated with unresolved optical sources, i.e. they are most likely distant QSOs or Galactic stars.

Figure 10 shows that the distributions of the RA and Dec angular separations between the positions of X-ray sources and their optical counterparts ($LR > 1.5$) can be approximated with Gaussians with FWHM of about 1 arcsec. The overall astrometric accuracy of the X-ray catalogue is therefore about 1.5 arcsec. This justifies the use of a normal distribution with a standard deviation of 1.5 arcsec for $f(r)$ in equation 9.

The XMM/SDSS source catalogue is made available to the public¹ in fits format. The information included in that file is listed in Appendix B.

10 THE COLOUR-MAGNITUDE DIAGRAM

One of the main motivations for the XMM/SDSS serendipitous survey is the compilation of a low redshift X-ray selected AGN sam-

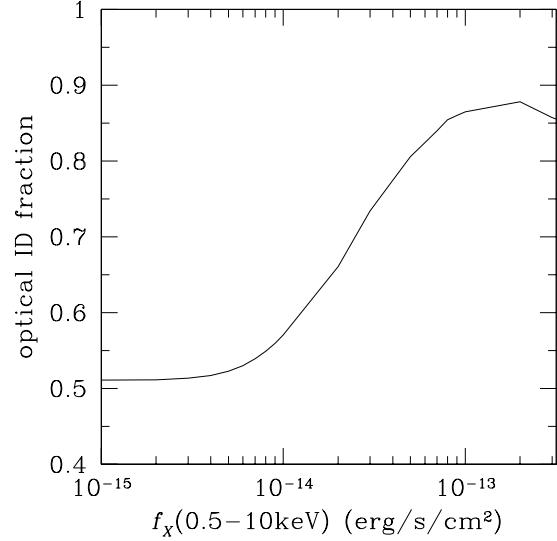


Figure 6. Distribution of the fraction of full-band selected X-ray point sources with optical counterparts as a function of 0.5–10 keV flux limit.

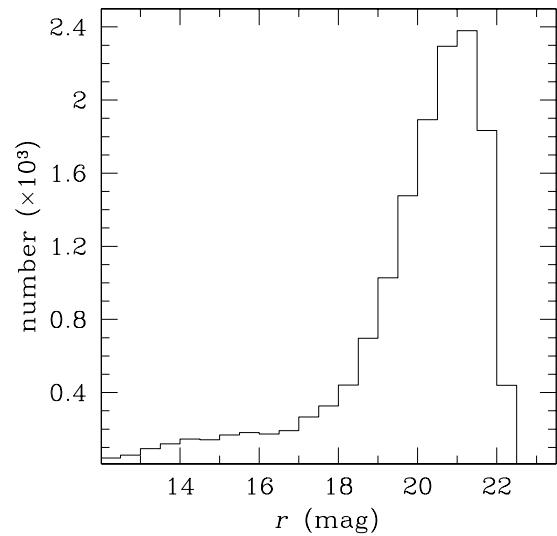


Figure 7. SDSS r -band magnitude distribution of full band selected X-ray sources in the XMM/SDSS survey.

ple which can be compared with high redshift AGN identified in deep X-ray surveys. In this section we present preliminary results on the comparison of the colour-magnitude diagram (CMD) of X-ray AGN at $z \approx 0.1$ and $z \approx 0.8$.

The low redshift X-ray sample consists of 293 hard-band (2–8 keV) detections in the XMM/SDSS survey with $0.03 < z < 0.2$ and $r < 17.77$ mag after correcting for Galactic extinction (Schlegel et al. 1998). The magnitude cut corresponds to the limit of the SDSS Main Galaxy Sample ($r < 17.77$ mag Strauss et al. 2002), which provides the vast majority of redshifts in the SDSS. The photometry is from the New York University Value-Added Galaxy Catalog (NYU-VAGC Blanton et al. 2005) which corresponds to the SDSS DR7 (Abazajian et al. 2009). This is because

¹ www.astro.noa.gr/~age

Table 1. X-ray point source catalogue details

Sample	total number of sources	$LR > 1.5$	spectroscopy	$r < 17.77$ mag	$r < 17.77$ mag AND spec-z	$r < 17.77$ mag AND $0.03 \leq z_{spec} \leq 0.2$
(1)	(2)	(3)	(4)	(5)	(6)	(7)
full band	35727	18246	1949	1920	844	560
soft band	33162	17631	1943	2018	840	552
hard band	14733	8803	1237	728	496	298
vhrd band	2245	1722	544	375	297	174
uhrd band	258	232	143	140	111	58

The columns are: (1) X-ray sample; (2) total number of X-ray point sources; (3) sources with optical counterpart ($LR > 1.5$); (4) X-ray sources with $LR > 1.5$ and SDSS spectroscopy (5) X-ray sources with optical counterparts brighter than $r = 17.77$ mag (petrosian magnitude, AB system) after correcting for Galactic extinction; (6) X-ray sources with $r \leq 17.77$ mag and spectroscopic redshift measurement; (7) X-ray sources with $r \leq 17.77$ mag and redshifts in the range $0.03 - 0.2$, i.e. the selection criteria used to construct the colour magnitude diagram of section 9. Note however, that in section 9 we use the photometry from the New York University Value-Added Galaxy Catalog (Blanton et al. 2005) and therefore there are small differences compared in the sample size compared to the numbers listed in column (7).

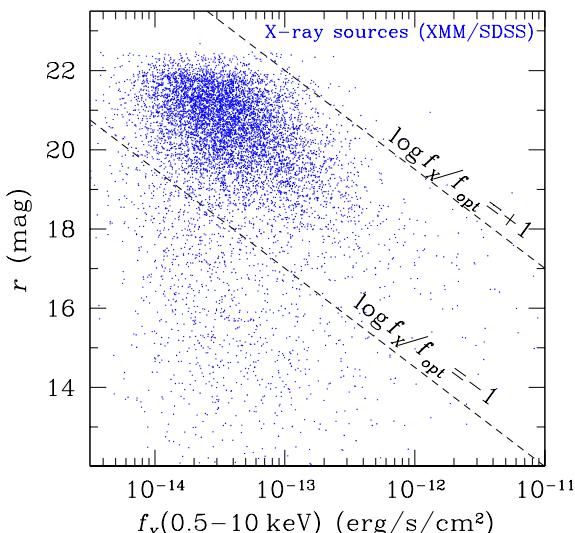


Figure 8. r -band magnitude against 0.5–10 keV flux for the XMM/SDSS X-ray sources with optical counterparts. The diagonal lines correspond to $\log f_X/f_{opt} = \pm 1$.

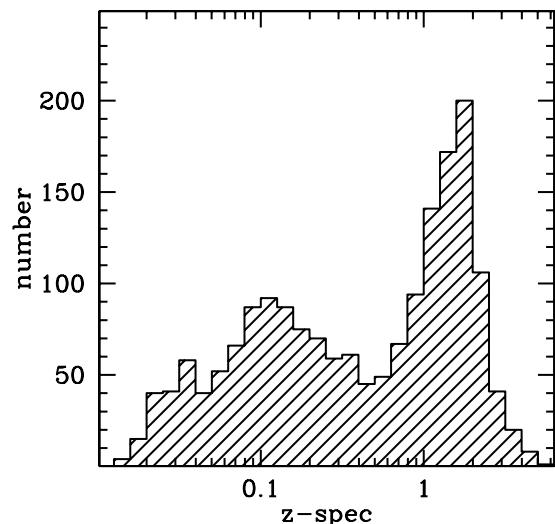


Figure 9. Redshift distribution of the XMM/SDSS X-ray sources in logarithmic bins of size $\Delta \log z = 0.1$.

the NYU-VAGC provides better photometric calibration of the SDSS data (Padmanabhan et al. 2008) compared to the DR7 and an accurate description of the SDSS window function. A small number of XMM/SDSS sources with $0.03 < z < 0.2$ and $r < 17.77$ mag (total of 5, see column 7 of Table 1) do not have photometric measurements in the NYU-VAGC and have been excluded from the final sample. Visual inspection shows that all of them are close to the edge of the SDSS field of view. The choice of the 2–8 keV X-ray band is because it is less affected by obscuration biases compared to softer energies and the XMM’s sensitivity remains high in that band thereby resulting in a sample that is sufficiently large for statistical studies. The CMD of the overall galaxy population at $z \approx 0.1$ is defined by selecting a total of 13,619 NYU-VAGC sources with $r < 17.77$ mag that overlap with the XMM/SDSS survey fields and have spectroscopic redshifts in the interval $0.03 < z < 0.2$.

The KCORRECT version 4.2 routines developed by Blanton & Roweis (2007) are used to fit spectral models to the NYU-VAGC $ugriz$ photometry and then estimate rest-frame magnitudes (AB system) in the $0.1u$, $0.1g$ bands, which are the SDSS u , g filters shifted to $z = 0.1$. The advantage of this choice of bandpasses is twofold. They are close to rest-frame for our $z \approx 0.1$ XMM/SDSS

sample and also have effective wavelengths that are similar to the rest-frame effective wavelengths of the DEEP2 R and I filters at $z = 1$ (Blanton 2006).

The high redshift CMD is constructed using data from the AEGIS (Davis et al. 2007). The spectroscopic sample consists of 6797 galaxies observed as part of the DEEP2 survey (Data Release 3) with secure redshift determinations ($> 90\%$ confidence level, quality flag $Q \geq 3$; Davis et al. 2007) in the interval $0.6 < z < 1.2$. The X-ray data in the AEGIS are from the Chandra survey of the Extended Groth Strip (AEGIS-X survey, Laird et al. 2009). The X-ray sources are identified with DEEP2 galaxies using the Likelihood Ratio methodology described in Laird et al. (2009). The spectroscopy for X-ray sources is from a variety of sources: the DEEP2 redshift survey, the SDSS and a number of spectroscopic programs that have targeted galaxies in the original Groth Strip (Weiner et al. 2005). We select a total of 115 X-ray sources in the spectroscopic redshift interval $0.6 < z < 1.2$ detected in the 0.5–7 keV band of the AEGIS-X survey. The median redshift of the sample is 0.8. The choice of the X-ray band for the detection of sources is to ensure that our $z \approx 0.8$ sample is selected at similar rest-frame energies (1–14 keV) as the $z \approx 0.1$ sample (2–8 keV), thereby facilitating the comparison. For the estimation of rest-frame quantities we use the

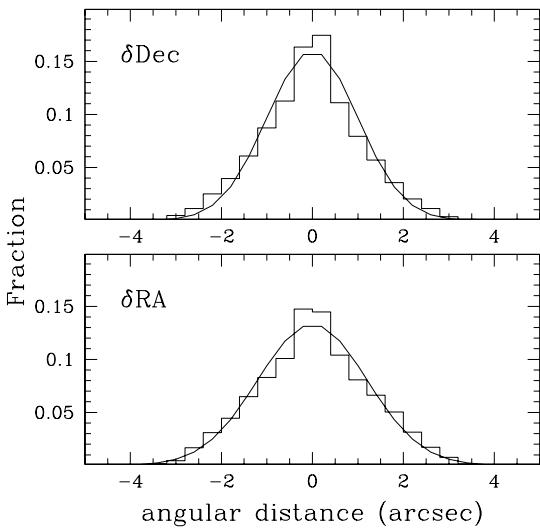


Figure 10. Angular separation distribution between the X-ray sources and their optical counterparts in the SDSS ($LR > 1.5$). The top and bottom panels show the Dec and RA offset histograms respectively. The curves are Gaussians with FWHM of about 1 arcsec and provide a fair approximation of the observed distributions.

BRI photometry obtained as part of the DEEP2 survey (Coil et al. 2004). The KCORRECT routines are used to fit models to the *BRI* data and then estimate the rest-frame $^{0.1}u$ and $^{0.1}g$ -band magnitudes from the observed DEEP2 *R* and *I*-band photometry respectively.

Figure 11 compares the absorption corrected 2-10 keV X-ray luminosity and the hydrogen column density (N_H) distributions of X-ray sources in the AEGIS sample and the low redshift subset of the XMM/SDSS survey. The N_H of individual sources are determined from the hardness ratios between the soft (0.5-2 keV for both AEGIS and XMM/SDSS) and the hard (2-7 keV for AEGIS, 2-8 keV for XMM) X-ray bands assuming an intrinsic power-law X-ray spectrum with index $\Gamma = 1.9$ (e.g. Nandra & Pounds 1994). The derived column densities are then used to correct the X-ray luminosities for photoelectric absorption. The XMM/SDSS sample at $z \approx 0.1$ is sensitive to sources with $\log L_X(2 - 10 \text{ keV}) > 41.5 \text{ erg s}^{-1}$. In contrast the AEGIS detects X-ray sources brighter than $\log L_X(2 - 10 \text{ keV}) > 41.5 \text{ erg s}^{-1}$ at $z \approx 0.8$. To avoid luminosity dependent biases in the comparison of the CMDs of X-ray sources in the AEGIS and the low redshift subset of the XMM/SDSS we limit both samples to $\log L_X(2 - 10 \text{ keV}) > 41.5 \text{ erg s}^{-1}$. X-ray sources with luminosities $\lesssim 10^{42} \text{ erg s}^{-1}$ are often excluded from AGN studies because normal galaxies are believed to be an important source of contamination below this limit. Although normal galaxy *candidates* are present at faint X-ray luminosities, AGN remain the dominant population. In Georgakakis et al. (2007) for example, X-ray detected starburst galaxy candidates represent only about 20 per cent of the GOODS-North X-ray source population in the luminosity interval $10^{40} - 10^{42} \text{ erg s}^{-1}$. This fraction is also likely to be an upper limit as some of the galaxy candidates may turn out to be low luminosity AGN. Furthermore, Georgakakis (2008) cross-correlated the 2XMMp catalogue (Watson et al. 2009) with the SDSS spectroscopic database to select normal galaxies at $z \approx 0.1$. Such sources were found to represent 20 per cent of the overall X-ray population for luminosities

$10^{40} - 10^{42} \text{ erg s}^{-1}$. This fraction is also likely to be a strong function of luminosity within that interval. Contamination by non-AGN is therefore not a concern for the XMM/SDSS and the AEGIS X-ray source samples with $\log L_X(2 - 10 \text{ keV}) > 41.5 \text{ erg s}^{-1}$. This cut reduces the total number of low redshift ($0.03 < z < 0.2$, $r < 17.77 \text{ mag}$) XMM/SDSS X-ray sources from 293 to 209. The sample includes a total of 16 XMM targets, i.e. non serendipitous X-ray sources, most of which (13/16) are type-1 AGN with broad optical emission lines. We choose to keep them in the sample as they are not affecting the main results and conclusions.

Figure 12 compares the $^{0.1}(u - g)$ vs $M_{0.1}g$ CMD of galaxies and X-ray sources in the AEGIS and the low redshift XMM/SDSS survey sample. X-ray sources which either show broad emission lines in their optical spectra and/or a prominent nuclear point source in their broadband optical images (SDSS or HST/ACS survey of the AEGIS) are marked in Figure 12. The optical continuum of those sources is significantly contaminated by AGN light. Their colours do not provide information on their host galaxies and are therefore excluded from further analysis. As expected the optical colours of these sources are very blue. The $^{0.1}(u - g)$ distribution of the two samples are shown in Figure 13, which better demonstrates the bimodality of the galaxy colours in both the AEGIS and the SDSS. We set the separation between red sequence and blue cloud galaxies at $^{0.1}(u - g) = 1.4$ for both the AEGIS at $z \approx 0.8$ and the SDSS at $z \approx 0.1$. Although the dividing line between red and blue galaxies also depends on absolute magnitude (e.g. Bell et al. 2004; Willmer et al. 2006), for our purposes it is sufficient to apply a simple colour cut (e.g. Blanton 2007). X-ray sources with colours contaminated by AGN light are not plotted in Figure 13. It is interesting that X-ray AGN in both the AEGIS and the low redshift subset of the SDSS/XMM survey have similar rest-frame colour distributions. The Kolmogorov-Smirnov test estimates a probability of 50 per cent that the two samples are drawn from the same parent population. Also, the fraction of X-ray AGN with $^{0.1}(u - g) > 1.4$ is 50 ± 4 in the low redshift subset of the XMM/SDSS and 56 ± 5 in the AEGIS, where the errors are estimated using binomial statistics.

Before interpreting those fractions however, the differences in the selection function of the two samples need to be accounted for. X-ray AGN in the AEGIS include, on average, a higher fraction of AGN with high hydrogen column density, $N_H > 10^{22} \text{ cm}^{-2}$ (see Figure 11 right panel). The AGN light of those obscured AGN is less likely to affect the host galaxy colours. We account for these effects using the simulations described in the appendix C to explore how the rest-frame colours of the AEGIS X-ray sources at $z \approx 0.8$ would appear in the low redshift ($0.03 < z < 0.2$) subsample of the XMM/SDSS survey. In brief, AEGIS X-ray sources are assigned a random redshift in the range $0.03 - 0.2$. Their X-ray flux (2-10 keV band) and apparent magnitude (SDSS *r*-band) are estimated at this new redshift to determine if they fulfill the selection criteria (e.g. X-ray sensitivity curve, $r < 17.77 \text{ mag}$) of the low redshift subsample of the XMM/SDSS survey. We can then explore how the various redshift dependent selection effects modify the derived rest-frame colour distribution and if they alter the fraction of AGN in red hosts. The simulated distribution of $^{0.1}(u - g)$ colours is plotted in Figure 14. It is found that the fraction of X-ray sources with $^{0.1}(u - g) > 1.4$ is 55 per cent, i.e. the same as for the ≈ 0.8 X-ray sources. Differential selection effects are not important for the comparison of the CMD of X-ray sources in the AEGIS and the XMM/SDSS survey. We therefore conclude that the data are consistent with no evolution of the rest-frame colours of X-ray AGN from $z = 0.8$ to $z = 0.1$.

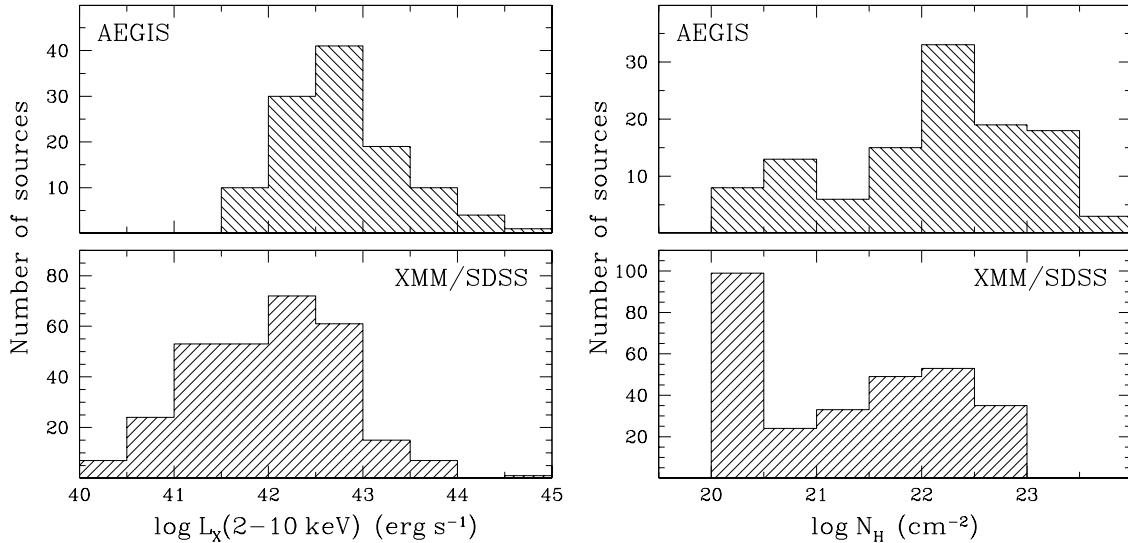


Figure 11. Left panel: 2-10 keV X-ray luminosity distribution of X-ray sources in the AEGIS (top) and the low redshift subset of the XMM/SDSS (bottom) samples. The X-ray luminosities are corrected for photoelectric absorption. In the case of AEGIS the 2-10 keV luminosity is estimated from the observed 0.5-7 keV count rate. The XMM/SDSS sample at $0.03 < z < 0.2$ reaches luminosities as faint as $10^{40} \text{ erg s}^{-1}$, while AEGIS at $0.6 < z < 1.2$ probes more luminous X-ray sources. **Right panel:** Histograms of the absorbing hydrogen column density, N_H in the AEGIS (top) and $z \approx 0.1$ XMM/SDSS (bottom) samples. This is determined from the observed hardness ratio of individual sources. For the AEGIS the hardness ratio is estimated from the count rates in the Chandra 0.5-2 and 2-7 keV spectral bands. In the case of the XMM/SDSS sample the hardness ratio is defined between the XMM's 0.5-2 and 2-8 keV bands.

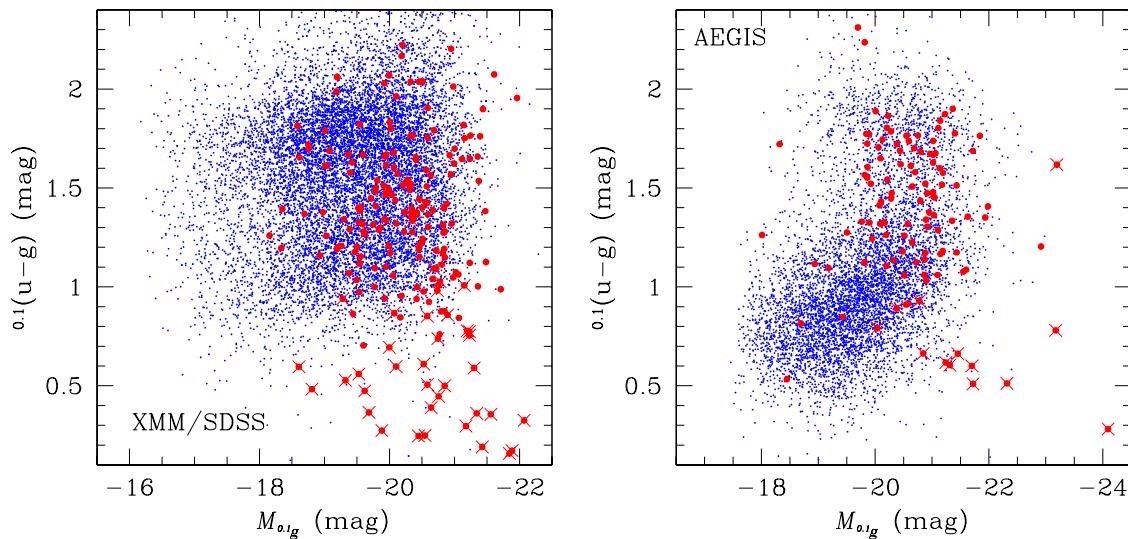


Figure 12. $0.1(u-g)$ vs $M_{0.1g}$. Colour magnitude diagram of galaxies (blue dots) and X-ray AGN (red circles) in the XMM/SDSS low redshift sample ($0.03 < z < 0.2$; left panel) and the AEGIS ($0.6 < z < 1.3$; right panel). Crosses indicate X-ray sources with rest-frame colours contaminated by AGN light.

11 DISCUSSION AND CONCLUSIONS

We present a new serendipitous X-ray survey, the XMM/SDSS, which is based on archival XMM observations and covers 122 deg^2 in the SDSS DR7 footprint. The size and sensitivity of this survey are well suited for low redshift X-ray AGN studies. The XMM/SDSS source catalogue is combined with the SDSS optical spectroscopy to compile one of the largest hard X-ray (2-8 keV) AGN samples to date in the redshift interval $0.03 < z < 0.2$. The

selection function of this sample is well defined and similar to that of deep, pencil-beam X-ray surveys, which probe AGN at moderate and high redshift. This greatly facilitates the study of the properties of AGN across redshift to shed light on the physical mechanism(s) responsible for the rapid decline of the accretion power of the Universe from $z \approx 1$ to the present day. Variations with redshift of the dominant SMBH fueling mode, which in some scenarios drive the evolution of the AGN population (e.g. Hopkins & Hernquist 2006;

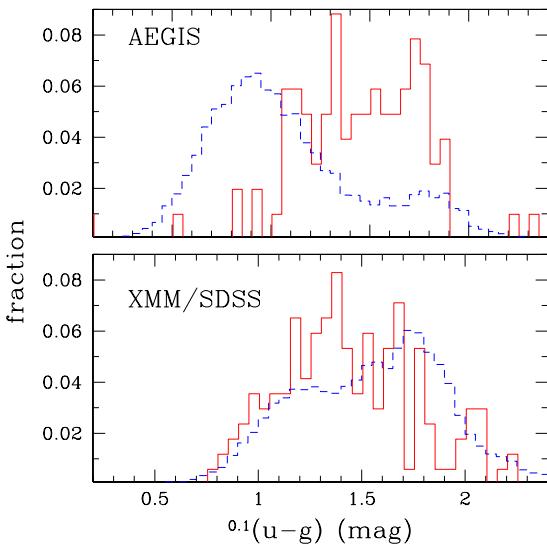


Figure 13. $0.1(u-g)$ colour distribution of galaxies (blue hatched histogram) and X-ray AGN (red solid histogram) in the XMM/SDSS low redshift sample ($0.03 < z < 0.2$; lower panel) and the AEGIS ($0.6 < z < 1.3$; upper panel). All histograms are normalised by the total number of sources in each sample.

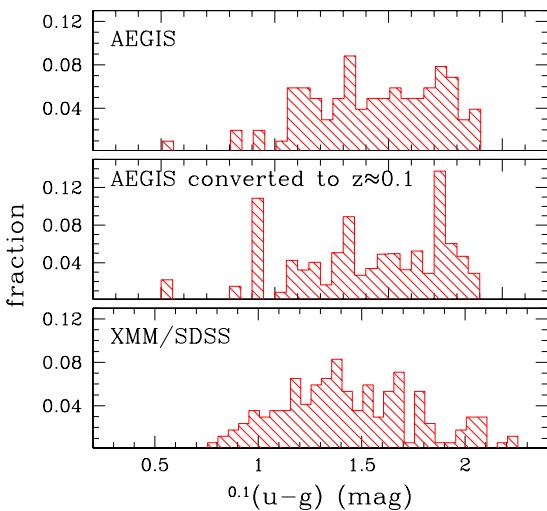


Figure 14. The top panel shows the $0.1(u-g)$ colour distribution of X-ray AGN at $0.6 < z < 1.2$ in the AEGIS field. The middle panel shows the colour distribution of that sample as it would appear in the redshift interval $0.03 < z < 0.2$ in our XMM/SDSS survey. For comparison the bottom plots the colour distribution of X-ray AGN in the range $0.03 < z < 0.2$ from the XMM/SDSS survey.

Hasinger 2008; Fanidakis et al. 2010), are expected to imprint on the properties of AGN, including those of their host galaxies. It has been proposed for example that major mergers dominate the AGN population at high redshift and bright luminosities, while stochastic accretion onto the SMBH of disk galaxies become important at low redshift and luminosities (Hopkins & Hernquist 2006). Alternatively, Fanidakis et al. (2010) argue that disk instabilities are the

main AGN fueling mode at early epochs, while hot gas accretion becomes increasingly more important at later times.

Systematic variations with redshift of the rest-frame colours of AGN hosts provide a first order test to the scenarios above. In the Fanidakis et al. (2010) prescription for example, the AGN population should shift to redder hosts with decreasing redshift. The opposite trend might be expected if disk instabilities dominate the accretion power at later epochs. Although a number of recent studies explored the position of AGN hosts on the colour-magnitude diagram (e.g. Hickox et al. 2009; Cardamone et al. 2010; Xue et al. 2010), none of them has investigated possible evolutionary effects with redshift, as it is done in this paper. We find no evidence for evolution of the CMD of X-ray AGN hosts from $z \approx 0.8$ (AEGIS-X) to $z \approx 0.1$ (XMM/SDSS). Apart from changes which are consistent with the evolution of the overall galaxy population between those redshifts, X-ray AGN are found in similar host galaxies, in terms of integrated colours, at both $z = 0.8$ and $z = 0.1$. This finding does not support the scenario in which the dominant mode of accretion changes with redshift since $z = 0.8$.

It should be noted however, that a number of biases may affect this conclusion. Although we have removed sources with colours dominated by the central engine (i.e. broad optical emission lines, bright nuclear point sources), we cannot exclude the possibility of residual contamination by scattered AGN light. Nevertheless, Pierce et al. (2010) studied the HST colour profiles of AGN hosts in the AEGIS survey. They concluded that the light from the active nucleus does not have a strong effect on the integrated colours of the majority of the AGN in their sample. Recently, Cardamone et al. (2010) suggest that about 25 per cent of the red cloud X-ray AGN at $z \approx 1$ are dusty star-forming systems, not evolved galaxies with a dominant old stellar population component. Systematic variations with redshift of the fraction of dust-reddened AGN hosts would impact our results and conclusions.

A striking feature of Figure 12 is the large fraction of blue type-1 AGN among low redshift X-ray sources compared to the AEGIS-X survey. The low redshift XMM/SDSS subsample includes XMM targets, which in their majority are broad line AGN (13/16). Also, the SDSS spectroscopic targets include apart from galaxies, QSOs at both low and high redshift, ROSAT X-ray sources, FIRST radio galaxies. The SDSS therefore has a higher spectroscopic sampling rate for type-1 AGN compared to the DEEP2 spectroscopic survey of the AEGIS. We conclude that selection effects are primarily responsible for the difference in the relative fraction of type-1 AGN between the AEGIS-X and the low redshift subset of the XMM/SDSS survey.

12 ACKNOWLEDGMENTS

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APPENDIX A: POINT SPREAD FUNCTION

The XMM Science Simulator (SCISIM) version 4.0.4 was used to produce point sources with monochromatic flux of 2×10^{-5} photons/mm/s at 1, 2, 3.5, 6 and 8.5 keV, i.e. the mean energies of the soft, full, hard, uhrd and vrhd bands respectively. The simulated sources were placed in a grid of positions separated by 90 arcsec on the PN, MOS1 and MOS2 cameras. A denser grid, although desirable, would be computationally expensive. Also, no simulations were produced for positions that are closer than 21 arcsec radius from CCD gaps, as in those cases the PSF size and shape cannot be reliably determined. The total exposure time was set to 50 ks. Because of technical issues related to the SCISIM code, this was split into 5 separate 10 ks simulations, which were then merged into a single event file. The source photon flux and the integration time were chosen to produce sufficient number of counts on the EPIC cameras to determine the PSF properties. Each simulation consists of a single energy, position and detector. There are 324, 361 and 380 PSFs at each energy for the PN, MOS1 and MOS2 cameras respectively.

The SCISIM generated ODF level files which were processed with SAS to produce event files, images and exposure maps for each simulation. The source counts were fit with ellipses using the IRAF (Image Reduction and Analysis Facility) task ELLIPSE of the STSDAS package, which was designed to fit the isophotes of optical galaxies. The task reads 2-dimensional images and fits ellipses to the isophotes at predefined semi-major axis values. During the ellipse fitting process the centre of the source was kept fixed. The position angle, PA, and ellipticity, ϵ , were left free parameters up to the semi-major axis value where S/N constraints terminated the fitting. Beyond that radius the position angle and ellipticity were kept constant. Parts of the image with exposure time equal to zero (e.g. CCD gaps) were masked out during the fitting process. The output of the task is a table that includes, among other parameters, the best fit ellipticity and position angle at different semi-major axis values. We used those parameters to estimate the PSF profile by extracting counts within elliptical annuli. The fraction of the pixels of the annuli that fall either outside the EPIC field of view or within CCD gaps is accounted for in the calculations. The cumulative radial distribution of the corrected counts is constructed, which is then used to estimate by linear interpolation the semi major axis radius, ϵ and PA of the 60, 70 and 80 per cent EEF of the PSF at different monochromatic energies and positions on the EPIC CCDs.

APPENDIX B: TABLE COLUMNS

B1 source detection parameters

UXID (30A): Unique X-ray source identification string. The unique name of the each source is a string which is composed from the project name, followed by the number of the pointing on which the source is detected or zero in the case of projects with a single XMM observation and ending with a sequential number.

BOX_ID_SRC (J): identification number of a source on individual pointings or observations in the case of projects with a single XMM observation.

RA (D): Right Ascension in degrees of the X-ray source before correcting for systematic offsets.

DEC (D): Declination in degrees of the X-ray source before correcting for systematic offsets.

RADEC_ERR (E): Positional uncertainty in arcsec of the source

position. This parameter is estimated by the EWAVELET task of SAS.

CNT_TOTAL_FULL_70 (J): full band counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

BKG_TOTAL_FULL_70 (E): expected number of full band background counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

EXP_TOTAL_FULL_70 (E): full band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 70 per cent EEF ellipse.

PROB_FULL (E): Poisson probability that the source is spurious in the full band.

CNT_PN_FULL_70 (J): full band counts at the source position on the PN camera within the 70 per cent EEF ellipse.

BKG_PN_FULL_70 (E): expected number of full band background counts at the source position on the PN camera within the 70 per cent EEF ellipse.

EXP_PN_FULL_70 (E): PN camera exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M1_FULL_70 (J): full band counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

BKG_M1_FULL_70 (E): expected number of full band background counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

EXP_M1_FULL_70 (E): MOS1 camera exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M2_FULL_70 (J): full band counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

BKG_M2_FULL_70 (E): expected number of full band background counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

EXP_M2_FULL_70 (E): MOS2 camera exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

EML_DET_ML_FULL (E): EMLDETECT detection likelihood in the full band for all EPIC cameras.

EML_EXT_ML_FULL (E): Likelihood that the source is extended in the full band. This parameter is estimated by the EMLDETECT task of SAS for all EPIC cameras.

CNT_TOTAL_SOFT_70 (J): soft band counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

BKG_TOTAL_SOFT_70 (E): expected number of soft band background counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

EXP_TOTAL_SOFT_70 (E): soft band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 70 per cent EEF ellipse.

PROB_SOFT (E): Poisson probability that the source is spurious in the soft band.

CNT_PN_SOFT_70 (J): soft band counts at the source position on the PN camera within the 70 per cent EEF ellipse.

BKG_PN_SOFT_70 (E): expected number of soft band background counts at the source position on the PN camera within the 70 per cent EEF ellipse.

EXP_PN_SOFT_70 (E): PN camera soft-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M1_SOFT_70 (J): soft band counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

BKG_M1_SOFT_70 (E): expected number of soft band background

counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

EXP_M1_SOFT_70 (E): MOS1 camera soft-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M2_SOFT_70 (J): soft band counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

BKG_M2_SOFT_70 (E): expected number of soft band background counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

EXP_M2_SOFT_70 (E): MOS2 camera soft-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

EML_DET_ML_SOFT (E): EMLDETECT detection likelihood in the soft band for all EPIC cameras.

EML_EXT_ML_SOFT (E): Likelihood that the source is extended in the soft band. This parameter is estimated by the EMLDETECT task of SAS for all EPIC cameras.

CNT_TOTAL_HARD_70 (J): hard band counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

BKG_TOTAL_HARD_70 (E): expected number of hard band background counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

EXP_TOTAL_HARD_70 (E): hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 70 per cent EEF ellipse.

PROB_HARD (E): Poisson probability that the source is spurious in the hard band.

CNT_PN_HARD_70 (J): hard band counts at the source position on the PN camera within the 70 per cent EEF ellipse.

BKG_PN_HARD_70 (E): expected number of hard band background counts at the source position on the PN camera within the 70 per cent EEF ellipse.

EXP_PN_HARD_70 (E): PN camera hard-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M1_HARD_70 (J): hard band counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

BKG_M1_HARD_70 (E): expected number of hard band background counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

EXP_M1_HARD_70 (E): MOS1 camera hard-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M2_HARD_70 (J): hard band counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

BKG_M2_HARD_70 (E): expected number of hard band background counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

EXP_M2_HARD_70 (E): MOS2 camera hard-band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

EML_DET_ML_HARD (E): EMLDETECT detection likelihood in the hard band for all EPIC cameras.

EML_EXT_ML_HARD (E): Likelihood that the source is extended in the hard band. This parameter is estimated by the EMLDETECT task of SAS for all EPIC cameras.

CNT_TOTAL_VHRD_70 (J): very hard band counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

BKG_TOTAL_VHRD_70 (E): expected number of very hard band background counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

EXP_TOTAL_VHRD_70 (E): very hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 70 per cent EEF ellipse.

PROB_VHRD (E): Poisson probability that the source is spurious in the very hard band.

CNT_PN_VHRD_70 (J): very hard band counts at the source position on the PN camera within the 70 per cent EEF ellipse.

BKG_PN_VHRD_70 (E): expected number of very hard band background counts at the source position on the PN camera within the 70 per cent EEF ellipse.

EXP_PN_VHRD_70 (E): PN camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M1_VHRD_70 (J): very hard band counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

BKG_M1_VHRD_70 (E): expected number of very hard band background counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

EXP_M1_VHRD_70 (E): MOS1 camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M2_VHRD_70 (J): very hard band counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

BKG_M2_VHRD_70 (E): expected number of very hard band background counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

EXP_M2_VHRD_70 (E): MOS2 camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

EML_DET_ML_VHRD (E): EMLDETECT detection likelihood in the very hard band for all EPIC cameras.

EML_EXT_ML_VHRD (E): Likelihood that the source is extended in the very hard band. This parameter is estimated by the EMLDETECT task of SAS for all EPIC cameras.

CNT_TOTAL_UHRD_70 (J): ultra hard band counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

BKG_TOTAL_UHRD_70 (E): expected number of ultra hard band background counts at the source position from all EPIC cameras within the 70 per cent EEF ellipse.

EXP_TOTAL_UHRD_70 (E): ultra hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 70 per cent EEF ellipse.

PROB_UHRD (E): Poisson probability that the source is spurious in the ultra hard band.

CNT_PN_UHRD_70 (J): ultra hard band counts at the source position on the PN camera within the 70 per cent EEF ellipse.

BKG_PN_UHRD_70 (E): expected number of ultra hard band background counts at the source position on the PN camera within the 70 per cent EEF ellipse.

EXP_PN_UHRD_70 (E): PN camera ultra hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M1_UHRD_70 (J): ultra hard band counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

BKG_M1_UHRD_70 (E): expected number of ultra hard band background counts at the source position on the MOS1 camera within the 70 per cent EEF ellipse.

EXP_M1_UHRD_70 (E): MOS1 camera ultra hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

CNT_M2_UHRD_70 (J): ultra hard band counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

BKG_M2_UHRD_70 (E): expected number of ultra hard band background counts at the source position on the MOS2 camera within the 70 per cent EEF ellipse.

EXP_M2_UHRD_70 (E): MOS2 camera ultra hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 70 per cent EEF ellipse.

EML_DET_ML_UHRD (E): EMLDETECT detection likelihood in the ultra hard band for all EPIC cameras.

EML_EXT_ML_UHRD (E): Likelihood that the source is extended in the ultra hard band. This parameter is estimated by the EMLDETECT task of SAS for all EPIC cameras.

RA_CORR (D): X-ray source Right Ascension in degrees corrected for systematic offsets of the XMM pointing.

DEC_CORR (D): X-ray source Declination in degrees corrected for systematic offsets of the XMM pointing.

LII_CORR (D): X-ray source Galactic longitude in degrees corresponding to RA_CORR and DEC_CORR

BIL_CORR (D) X-ray source Galactic latitude in degrees corresponding to RA_CORR and DEC_CORR

B2 Source parameters related to flux estimation

CNT_TOTAL_FULL_80 (J): full band counts at the source position from all EPIC cameras estimated within the 80 per cent EEF ellipse.

BKG_TOTAL_FULL_80 (E): expected number of full band background counts at the source position from all EPIC cameras within the 80 per cent EEF ellipse.

EXP_TOTAL_FULL_80 (E): full band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_PN_FULL_80 (J): full band counts at the source position on the PN camera within the 80 per cent EEF ellipse.

BKG_PN_FULL_80 (E): expected number of full band background counts at the source position on the PN camera within the 80 per cent EEF ellipse.

EXP_PN_FULL_80 (E): PN camera full band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M1_FULL_80 (J): full band counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

BKG_M1_FULL_80 (E): expected number of full band background counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

EXP_M1_FULL_80 (E): MOS1 camera full band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M2_FULL_80 (J): full band counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

BKG_M2_FULL_80 (E): expected number of full band background counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

EXP_M2_FULL_80 (E): MOS2 camera full band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_TOTAL_SOFT_80 (J): soft band counts at the source position from all EPIC cameras estimated within the 80 per cent EEF ellipse.

BKG_TOTAL_SOFT_80 (E): expected number of soft band background counts at the source position from all EPIC cameras within the 80 per cent EEF ellipse.

EXP_TOTAL_SOFT_80 (E): soft band exposure time in seconds at

the source position from all EPIC cameras. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_PN_SOFT_80 (J): soft band counts at the source position on the PN camera within the 80 per cent EEF ellipse.

BKG_PN_SOFT_80 (E): expected number of soft band background counts at the source position on the PN camera within the 80 per cent EEF ellipse.

EXP_PN_SOFT_80 (E): PN camera soft band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M1_SOFT_80 (J): soft band counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

BKG_M1_SOFT_80 (E): expected number of soft band background counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

EXP_M1_SOFT_80 (E): MOS1 camera soft band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M2_SOFT_80 (J): soft band counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

BKG_M2_SOFT_80 (E): expected number of soft band background counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

EXP_M2_SOFT_80 (E): MOS2 camera soft band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_TOTAL_HARD_80 (J): hard band counts at the source position from all EPIC cameras estimated within the 80 per cent EEF ellipse.

BKG_TOTAL_HARD_80 (E): expected number of hard band background counts at the source position from all EPIC cameras within the 80 per cent EEF ellipse.

EXP_TOTAL_HARD_80 (E): hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_PN_HARD_80 (J): hard band counts at the source position on the PN camera within the 80 per cent EEF ellipse.

BKG_PN_HARD_80 (E): expected number of hard band background counts at the source position on the PN camera within the 80 per cent EEF ellipse.

EXP_PN_HARD_80 (E): PN camera hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M1_HARD_80 (J): hard band counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

BKG_M1_HARD_80 (E): expected number of hard band background counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

EXP_M1_HARD_80 (E): MOS1 camera hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M2_HARD_80 (J): hard band counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

BKG_M2_HARD_80 (E): expected number of hard band background counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

EXP_M2_HARD_80 (E): MOS2 camera hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_TOTAL_VHRD_80 (J): very hard band counts at the source position from all EPIC cameras estimated within the 80 per cent EEF ellipse.

BKG_TOTAL_VHRD_80 (E): expected number of very hard band background counts at the source position from all EPIC cameras within the 80 per cent EEF ellipse.

EXP_TOTAL_VHRD_80 (E): very hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_PN_VHRD_80 (J): very hard band counts at the source position on the PN camera within the 80 per cent EEF ellipse.

BKG_PN_VHRD_80 (E): expected number of very hard band background counts at the source position on the PN camera within the 80 per cent EEF ellipse.

EXP_PN_VHRD_80 (E): PN camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M1_VHRD_80 (J): very hard band counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

BKG_M1_VHRD_80 (E): expected number of very hard band background counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

EXP_M1_VHRD_80 (E): MOS1 camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M2_VHRD_80 (J): very hard band counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

BKG_M2_VHRD_80 (E): expected number of very hard band background counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

EXP_M2_VHRD_80 (E): MOS2 camera very hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_TOTAL_UHRD_80 (J): ultra hard band counts at the source position from all EPIC cameras estimated within the 80 per cent EEF ellipse.

BKG_TOTAL_UHRD_80 (E): expected number of ultra hard band background counts at the source position from all EPIC cameras within the 80 per cent EEF ellipse.

EXP_TOTAL_UHRD_80 (E): ultra hard band exposure time in seconds at the source position from all EPIC cameras. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_PN_UHRD_80 (J): ultra hard band counts at the source position on the PN camera within the 80 per cent EEF ellipse.

BKG_PN_UHRD_80 (E): expected number of ultra hard band background counts at the source position on the PN camera within the 80 per cent EEF ellipse.

EXP_PN_UHRD_80 (E): PN camera ultra hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M1_UHRD_80 (J): ultra hard band counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

BKG_M1_UHRD_80 (E): expected number of ultra hard band background counts at the source position on the MOS1 camera within the 80 per cent EEF ellipse.

EXP_M1_UHRD_80 (E): MOS1 camera ultra hard band exposure time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

CNT_M2_UHRD_80 (J): ultra hard band counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

BKG_M2_UHRD_80 (E): expected number of ultra hard band background counts at the source position on the MOS2 camera within the 80 per cent EEF ellipse.

EXP_M2_UHRD_80 (E): MOS2 camera ultra hard band exposure

time in seconds at the source position. This parameter is the average exposure time within the 80 per cent EEF ellipse.

ECF_PN_FULL (E): PN full band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M1_FULL (E): MOS1 full band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M2_FULL (E): MOS2 full band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_PN_SOFT (E): PN soft band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M1_SOFT (E): MOS1 soft band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M2_SOFT (E): MOS2 soft band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_PN_HARD (E): PN hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M1_HARD (E): MOS1 hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M2_HARD (E): MOS2 hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_PN_VHRD (E): PN very hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M1_VHRD (E): MOS1 very hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M2_VHRD (E): MOS2 very hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_PN_UHRD (E): PN ultra hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M1_UHRD (E): MOS1 ultra hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

ECF_M2_UHRD (E): MOS2 ultra hard band counts to flux conversion factor in units of 10^{11} counts/($\text{erg s}^{-1} \text{cm}^{-2}$).

HR_PN (E): PN hardness ratio. A value of -9.99 for HR_PN, HR_PN_UPPER and HR_PN_LOWER (for the definition of the latter two parameters see below) means that the HR_PN could not be estimated, e.g. the source lies outside the PN FOV.

HR_PN_UPPER (E): 68 per cent confidence level upper limit in HR_PN.

HR_PN_LOWER (E): 68 per cent confidence level lower limit in HR_PN.

HR_M1 (E): MOS1 hardness ratio. A value of -9.99 for HR_M1, HR_M1_UPPER and HR_M1_LOWER means that the HR_M1 could not be estimated, e.g. the source lies outside the MOS1 FOV.

HR_M1_UPPER (E): 68 per cent confidence level upper limit in HR_M1.

HR_M1_LOWER (E): 68 per cent confidence level lower limit in HR_M1.

HR_M2 (E): MOS2 hardness ratio. A value of -9.99 for HR_M2, HR_M2_UPPER and HR_M2_LOWER means that the HR_M2 could not be estimated, e.g. the source lies outside the MOS2 FOV.

HR_M2_UPPER (E): 68 per cent confidence level upper limit in HR_M2.

HR_M2_LOWER (E): 68 per cent confidence level lower limit in HR_M2.

FLUX_FULL (E): 0.5-10 keV flux in $\text{erg s}^{-1} \text{cm}^{-2}$. In the case of upper limit in the flux this parameter is set to zero.

FLUX_FULL_LOWER (E): 0.5-10 keV flux lower limit uncertainty (68 per cent). In the case of upper limit in the flux this parameter is set to zero.

FLUX_FULL_UPPER (E): 0.5-10 keV flux upper limit (68 per cent)

FLUX_SOFT (E): 0.5-2 keV flux in $\text{erg s}^{-1} \text{cm}^{-2}$.

FLUX_SOFT_LOWER (E): 0.5-2 keV flux lower limit (68 per cent)
 FLUX_SOFT_UPPER (E): 0.5-2 keV flux upper limit (68 per cent)
 FLUX_HARD (E): 2-10 keV flux in $\text{erg s}^{-1} \text{cm}^{-2}$.
 FLUX_HARD_LOWER (E): 2-10 keV flux lower limit (68 per cent)
 FLUX_HARD_UPPER (E): 2-10 keV flux upper limit (68 per cent)
 FLUX_VHRD (E): 5-10 keV flux in $\text{erg s}^{-1} \text{cm}^{-2}$.
 FLUX_VHRD_LOWER (E): 5-10 keV flux lower limit (68 per cent)
 FLUX_VHRD_UPPER (E): 5-10 keV flux upper limit (68 per cent)
 FLUX_UHRD (E): 7.5-12 keV flux in $\text{erg s}^{-1} \text{cm}^{-2}$.
 FLUX_UHRD_LOWER (E): 7.5-12 keV flux lower limit (68 per cent)
 FLUX_UHRD_UPPER (E): 7.5-12 keV flux upper limit (68 per cent)
 LG_NH_GAL (E): Base 10 logarithm of the Galactic hydrogen column density in the direction of the source in units of cm^{-3} estimated from the HI map of Kalberla et al. (2005).
 FILTER_PN (10A): XMM's PN filter.
 FILTER_MOS1 (10A): XMM's MOS1 filter.
 FILTER_MOS1 (10A): XMM's MOS2 filter.

B3 X-ray optical identification parameters

SDSS_ID (20A): SDSS source identification number.
 LR_SOFT (E): Likelihood of the optical counterpart for X-ray sources in the soft-band selected sample.
 LR_HARD (E): Likelihood of the optical counterpart for X-ray sources in the hard-band selected sample.
 LR_VHRD (E): Likelihood of the optical counterpart for X-ray sources in the very hard band selected sample.
 LR_UHRD (E): Likelihood of the optical counterpart for X-ray sources in the ultra hard band selected sample. sample.
 LR_FULL (E): Likelihood of the optical counterpart for X-ray sources in the full band selected sample.
 LR (E): Likelihood of the optical counterpart. This parameter is estimated for all X-ray sources independent of the detection band.
 REL_SOFT (E): Reliability of the optical counterpart for X-ray sources in the soft band selected sample.
 REL_HARD (E): Reliability of the optical counterpart for X-ray sources in the hard band selected sample.
 REL_VHRD (E): Reliability of the optical counterpart for X-ray sources in the very hard band selected sample.
 REL_UHRD (E): Reliability of the optical counterpart for X-ray sources in the ultra hard -band selected sample.
 REL_FULL (E): Reliability of the optical counterpart for X-ray sources in the full band selected sample.
 REL (E): Reliability of the optical counterpart. This parameter is estimated for all X-ray sources independent of the detection band.
 DR_ARCSEC (E): angular separation in arcsec between the X-ray source and the optical counterpart.
 DR_RA_ARCSEC (E): RA offset in arcsec between the X-ray source and the optical counterpart.
 DR_DEC_ARCSEC (E): DEC offset in arcsec between the X-ray source and the optical counterpart

B4 SDSS photometry parameters

petro_u (E): SDSS *u*-band petrosian magnitude of the optical counterpart.

petro_g (E): SDSS *g*-band petrosian magnitude of the optical counterpart.
 petro_r (E): SDSS *r*-band petrosian magnitude of the optical counterpart.
 petro_i (E): SDSS *i*-band petrosian magnitude of the optical counterpart.
 petro_z (E): SDSS *z*-band petrosian magnitude of the optical counterpart.
 petro_err_u (E): SDSS *u*-band petrosian magnitude error of the optical counterpart.
 petro_err_g (E): SDSS *g*-band petrosian magnitude error of the optical counterpart.
 petro_err_r (E): SDSS *r*-band petrosian magnitude error of the optical counterpart.
 petro_err_i (E): SDSS *i*-band petrosian magnitude error of the optical counterpart.
 petro_err_z (E): SDSS *z*-band petrosian magnitude error of the optical counterpart.
 model_u (E): SDSS *u*-band model magnitude of the optical counterpart.
 model_g (E): SDSS *g*-band model magnitude of the optical counterpart.
 model_r (E): SDSS *r*-band model magnitude of the optical counterpart.
 model_i (E): SDSS *i*-band model magnitude of the optical counterpart.
 model_z (E): SDSS *z*-band model magnitude of the optical counterpart.
 model_err_u (E): SDSS *u*-band model magnitude error of the optical counterpart.
 model_err_g (E): SDSS *g*-band model magnitude error of the optical counterpart.
 model_err_r (E): SDSS *r*-band model magnitude error of the optical counterpart.
 model_err_i (E): SDSS *i*-band model magnitude error of the optical counterpart.
 model_err_z (E): SDSS *z*-band model magnitude error of the optical counterpart.
 psf_u (E): SDSS *u*-band psf magnitude of the optical counterpart
 psf_g (E): SDSS *g*-band psf magnitude of the optical counterpart
 psf_r (E): SDSS *r*-band psf magnitude of the optical counterpart
 psf_i (E): SDSS *i*-band psf magnitude of the optical counterpart
 psf_z (E): SDSS *z*-band psf magnitude of the optical counterpart.
 psf_err_u (E): SDSS *u*-band psf magnitude error of the optical counterpart.
 psf_err_g (E): SDSS *u*-band psf magnitude error of the optical counterpart.
 psf_err_r (E): SDSS *u*-band psf magnitude error of the optical counterpart.
 psf_err_i (E): SDSS *u*-band psf magnitude error of the optical counterpart.
 psf_err_z (E): SDSS *u*-band psf magnitude error of the optical counterpart.
 ext_u (E): Galactic extinction in magnitudes in the SDSS *u*-band.
 ext_g (E): Galactic extinction in magnitudes in the SDSS *g*-band.
 ext_r (E): Galactic extinction in magnitudes in the SDSS *r*-band.
 ext_i (E): Galactic extinction in magnitudes in the SDSS *i*-band.
 ext_z (E): Galactic extinction in magnitudes in the SDSS *z*-band.
 p_type (J): SDSS photometric type, 3 for optically extended source, 6 for stellar object.
 GAL_LAT (D): Galactic latitude of the counterpart.
 GAL_LONG (D): Galactic longitude of the counterpart.

B5 SDSS photometric and spectroscopic redshift parameters

SDSS_SPEC_ID (20A): SDSS spectroscopic identification number, if available.

z (D): spectroscopic redshift

ztype(D): spectroscopic redshift type

z_phot1 (D): Neural Network photometric redshift estimated by the CC2 method described by Oyaizu et al. (2008).

z_err_1 (D): photometric redshift error for the CC2 method of Oyaizu et al. (2008).

z_phot2 (D): Neural Network photometric redshift estimated by the D1 method described by Oyaizu et al. (2008).

z_err_2 (D): photometric redshift error for the D1 method of Oyaizu et al. (2008).

APPENDIX C: SIMULATIONS: THE CMD OF X-RAY SOURCES IN THE AEGIS AS WOULD APPEAR AT LOW REDSHIFT

Differential selection effects in the comparison of the CMD of X-ray AGN in the AEGIS sample and the low redshift subset of the XMM/SDSS survey are accounted for by performing simulations to explore how the rest-frame colours of the AEGIS X-ray sources at $z \approx 0.8$ would appear in our low redshift ($0.03 < z < 0.2$) XMM/SDSS survey. We choose to transform the high redshift sample to low redshifts. This is because the Chandra observations of the AEGIS field select AGN at rest-frame energies 1-14 keV at $z \approx 0.8$. That sample is therefore less affected by obscuration compared to the XMM/SDSS survey, which probes lower rest frame energies ($\approx 2 - 8$ keV) at $z < 0.2$. The simulations proceed as follows. AEGIS X-ray sources are randomly selected by weighting each one by $1/V_{max}$, the inverse of the maximum comoving volume within which it can be detected. Each source is placed at a random redshift in the range $0.03 - 0.2$ and its SDSS r -band magnitude is estimated using the KCORRECT routines. The source is retained in the sample if it is brighter than the magnitude limit $r = 17.77$ mag of the XMM/SDSS low redshift survey. The 2-10 keV X-ray flux, $f_X(2 - 10 \text{ keV})$, of the source is also determined from the X-ray luminosity, $L_X(0.5 - 10 \text{ keV})$, and hydrogen column density, N_H . The sensitivity curve of the XMM/SDSS survey is an estimate of the probability of detecting a source with given flux within the surveyed area. A random number is produced between 0 and the maximum value of the XMM/SDSS survey 2-10 keV sensitivity curve. If the random number is lower than the value of the sensitivity curve at flux $f_X(2 - 10 \text{ keV})$ the source is retained in the sample. A total of 10,000 realisations are performed.

The estimation of the maximum volume takes into account the optical magnitude limits for spectroscopy in the AEGIS, $18.5 < R < 24.1$ mag, the varying X-ray sensitivity in that field and the spectroscopic sampling rate. For source i at redshift z , with intrinsic X-ray luminosity L_X , obscuring hydrogen column density N_H and absolute optical magnitude $M_{0,1}$ the maximum volume is

$$V_{max,i}(L_X, N_H, M, z) = \frac{c}{H_0} \int_{z1}^{z2} \Omega(L_X, z) w_i \frac{dV}{dz} dz dL, \quad (\text{C1})$$

where dV/dz is the volume element per redshift interval dz and w_i is a weight that depends on the observed $B - R$, $R - I$ colours and the R -band magnitude of the X-ray sources and corrects for the spectroscopic sampling rate of the X-ray population. The weights are estimated using the “minimal” model described by Willmer et al. (2006). The integration limits are $z1 = \max(z_{bright}, z_L)$ and $z2 = \min(z_{faint}, z_U)$, where we have defined z_L, z_U the lower

and upper redshift limits of the sample and z_{faint}, z_{bright} the redshifts at which the source will become fainter than $R = 24.1$ mag and brighter than $R = 18.5$ mag, the optical magnitude limits for optical spectroscopy. $\Omega(L_X, z)$ is the solid angle of the X-ray survey available to a source with luminosity L_X and column density N_H at redshift z (corresponding to a flux f_X in the X-ray sensitivity curve). The estimation of the X-ray flux from L_X takes into account the intrinsic source N_H by adopting a power-law spectral energy distribution with $\Gamma = 1.9$ and photoelectric absorption cross sections as described by Morrison & McCammon (1983) for solar metallicity.